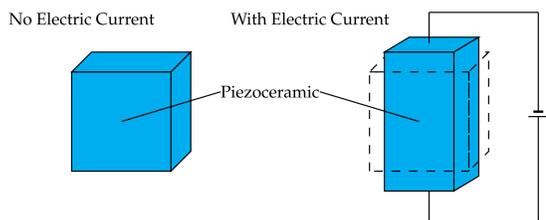


## Piezoelectric materials

Piezoelectric materials generate an electric charge when subjected to mechanical stress and also deform when an electric current is run through them.



Governing equations for piezoelectric materials:

$$\sigma_{ij,i} = 0, \quad D_{i,i} = 0$$

where  $\sigma$  is the stress and  $D$  is the electrical induction. The elastic and electric fields are related in the following manner [1]

$$\sigma_{ij} = C_{ijrs} \varepsilon_{rs} - e_{sji} E_s,$$

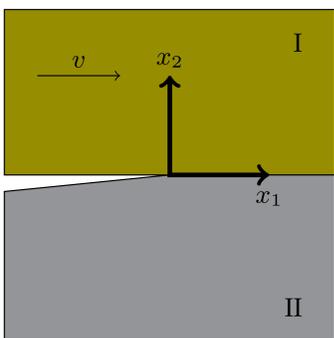
$$D_i = \omega_{is} E_s + e_{irs} \varepsilon_{rs},$$

where  $\varepsilon$  is strain,  $C$  is the material stiffness tensor,  $E$  is the electric field,  $\omega$  is permittivity and  $e$  is the piezoelectric tensor. An eigenvalue problem was then developed in [1] to find solutions in the form of the extended displacement and traction fields:

$$\mathbf{U} = (\mathbf{u}^T, \Phi)^T, \quad \mathbf{t} = (\sigma_{12}, \sigma_{22}, \sigma_{32}, D_2)^T.$$

## Problem formulation

Interfacial crack propagation in anisotropic bimaterials pose many interesting problems. The following geometry will be considered:



Perfect interface transmission conditions are considered:

- Continuation of displacements and tractions.
- Continuity of potential and charge for piezoelectric materials.

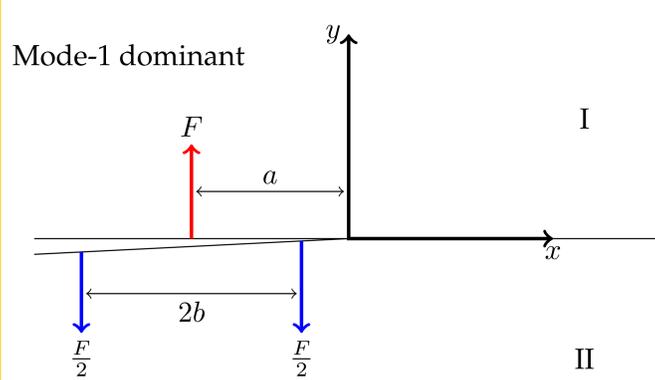
The following two examples are considered

1. Stress intensity factor analysis for a dynamic crack along the interface in an anisotropic bimaterial - no piezoelectricity [2]
2. Singular integral equations for a static interfacial crack between two piezoelectric materials [3]

## Acknowledgements

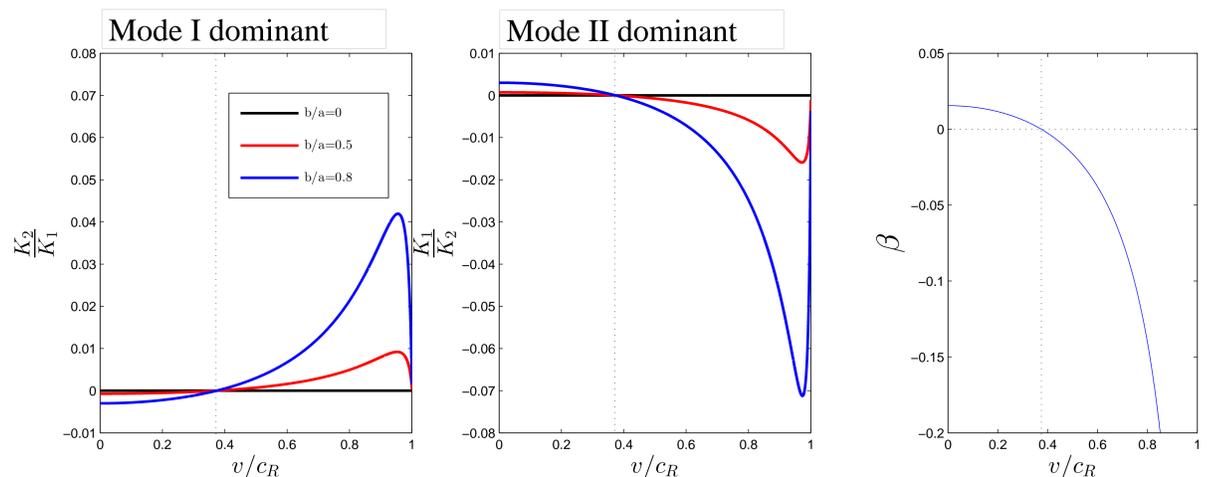
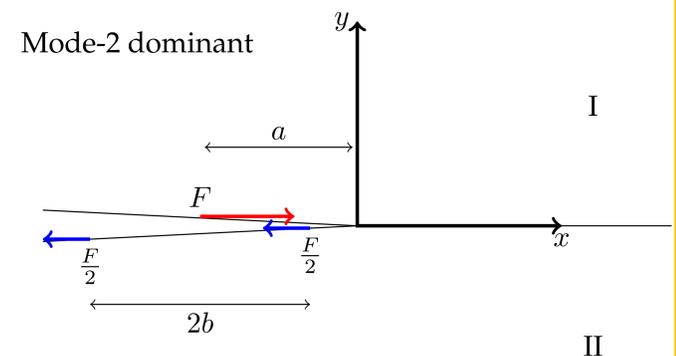
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## In-plane stress intensity factors for orthotropic bimaterials



- Both loading systems are symmetric when  $b = 0$ .
- Increased values of  $b$  leads to a higher degree of asymmetry.
- Finding the stress intensity factor,  $K = K_1 + iK_2$ , which plays a key role in the first order asymptotics of displacements and tractions at the crack tip.

- Orthotropic materials used.
- New co-ordinates:  $x = x_1 - vt, y = x_2$ .
- Steady-state formulation now possible.
- Limiting velocity - Rayleigh wave speed of most compliant material,  $c_R$ .
- Two asymmetric loading configurations considered.



- Symmetric loading ( $b/a = 0$ ) - stress intensity factor purely real for Mode-1 dominant loading or completely imaginary for Mode-2 dominant loading.
- Increased asymmetry leads to an increase in the non-dominant part of  $K$  - behaviour near  $c_R$  should be investigated practically.
- The Dundurs parameter,  $\beta$ , is real-valued and depends on the material parameters and crack velocity,  $v$ .
- There exists a particular velocity,  $v_0$ , where  $\beta(v_0) = 0$  - non-dominant part of  $K$  disappears at  $v_0$  regardless of loading asymmetry. Crack propagates straight along the interface. However,  $v_0$  does not exist for every bimaterial.

## Singular integral equations - out-of-plane piezoelectric effect

The example considered here has  $v = 0$  and poling in the  $x_3$  direction, therefore the in-plane fields are not affected by the piezoelectric effect and behave elastically. To study the piezo-effect, solutions are found in the form:  $\mathbf{u} = (u_3, \phi)^T$ ,  $\mathbf{t} = (\sigma_{32}, D_2)$ .

Implementing weight functions and the Betti formula leads to the following singular integral equations:

$$\langle \mathbf{p} \rangle(x_1) + \mathbf{Y} \llbracket \mathbf{p} \rrbracket(x_1) = -\frac{\mathbf{H}^{-1}}{\pi} \int_{-\infty}^0 \frac{1}{x_1 - \eta} \frac{\partial \llbracket \mathbf{u} \rrbracket^{(-)}}{\partial \eta} d\eta, \quad \text{for } x_1 < 0,$$

$$\mathbf{t}(x_1) = -\frac{\mathbf{H}^{-1}}{\pi} \int_{-\infty}^0 \frac{1}{x_1 - \eta} \frac{\partial \llbracket \mathbf{u} \rrbracket^{(-)}}{\partial \eta} d\eta, \quad \text{for } x_1 > 0,$$

where  $\mathbf{H}$  and  $\mathbf{Y}$  are bimaterial matrices depending on the material surface admittance tensors.  $\llbracket \mathbf{p} \rrbracket$  and  $\langle \mathbf{p} \rangle$  depend on the crack face loading,  $\mathbf{p}$ , in the following way:

$$\llbracket \mathbf{p} \rrbracket(x_1) = \mathbf{p}^+(x_1) - \mathbf{p}^-(x_1), \quad \langle \mathbf{p} \rangle(x_1) = \frac{1}{2} [\mathbf{p}^+(x_1) + \mathbf{p}^-(x_1)].$$

## References

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